A hierarchy of higher order and higher grade continua Application to the plasticity and fracture of metallic foams

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The micromorphic approach to plasticity

- Continuum thermomechanics
- Full micromorphic and microstrain theories
- Internal constraint in the micromorphic approach

2 Application to nickel foams

- Micromorphic anisotropic compressible plasticity model
- Simulation of crack propagation



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3 Micromechanical approach

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State space

• **observable and controllable variables** (temperature, strain...)

 $\{T, K_{\widetilde{n}}\}$

• **internal degrees of freedom** (controllable variables that account for some aspects of the microstructre)

 $\{\alpha, \nabla \alpha\}$

they have associated stresses and α or its associated force can be prescribed at the boundary

• **internal variables** are the remembrance of internal degrees of freedom; they cannot be controlled

 $\{\alpha\}$

gradient of internal variable (extra-entropy flux)

$$\{\alpha, \nabla \alpha\}$$

[Maugin, 1990] [Maugin, 1997]

The micromorphic approach (1)

• Start from an initial classical elastoviscoplastic model with internal variables

 $DOF0 = \{\underline{\mathbf{u}}\}, \quad STATE0 = \{\underline{\mathbf{F}}, \quad T, \quad \alpha\}$

 Select one variable φ ∈ STATE0 and introduce the associated micromorphic variable^χφ as an additional degree of freedom and, possibly, state variable:

$$DOF = \{\underline{\mathbf{u}}, \quad ^{\chi}\phi\}, \quad STATE = \{\underline{\mathbf{F}}, \quad T, \quad \alpha, \quad ^{\chi}\phi, \quad \nabla^{\chi}\phi\}$$

• Extend the power of internal forces

$$\mathcal{P}^{(i)}(\underline{\mathbf{v}}^{\star}, \overset{\times}{\psi} \dot{\phi}^{\star}) = -\int_{\mathcal{D}} p^{(i)}(\underline{\mathbf{v}}^{\star}, \overset{\times}{\psi} \dot{\phi}^{\star}) \, dV$$
$$p^{(i)}(\mathbf{v}^{\star}, \overset{\times}{\psi} \dot{\phi}^{\star}) = \boldsymbol{\sigma} : \boldsymbol{\nabla} \mathbf{v}^{\star} + \mathbf{a}^{\chi} \dot{\phi}^{\star} + \mathbf{b} \cdot \boldsymbol{\nabla}^{\chi} \dot{\phi}$$

a, **b** generalized stresses, *microforces* (Gurtin, 1996)

• Derive additional balance equation and boundary conditions

$$\operatorname{div} \underline{\mathbf{b}} - \mathbf{a} = \mathbf{0}, \quad \forall \underline{\mathbf{x}} \in \Omega, \quad \underline{\mathbf{b}} \, . \underline{\mathbf{n}} \, = \mathbf{a}^{c}, \forall \underline{\mathbf{x}} \, \in \partial \Omega$$

The micromorphic approach to plasticity

The micromorphic approach (2)

 More generally, in the presence of volume generalized forces: div (<u>**b**</u>-<u>**b**</u>^e)-a+a^e = 0, ∀<u>x</u> ∈ Ω, (<u>**b**</u>-<u>**b**</u>^e).<u>**n**</u> = a^c, ∀<u>x</u> ∈ ∂Ω

• Enhance the local balance of energy and the entropy inequality

$$\rho \dot{\epsilon} = p^{(i)} - \operatorname{div} \mathbf{\underline{q}} + \rho r, \quad -\rho(\dot{\psi} + \eta \dot{T}) + p^{(i)} - \frac{\mathbf{\underline{q}}}{T} \cdot \nabla T \ge 0$$

• Consider the constitutive functionals:

$$\psi = \hat{\psi}(\mathbf{F}^{e}, T, \alpha, ^{\chi}\phi, \nabla^{\chi}\phi), \ \eta = \hat{\eta}(\mathbf{F}^{e}, T, \alpha, ^{\chi}\phi, \nabla^{\chi}\phi)$$

$$\boldsymbol{\sigma} = \hat{\boldsymbol{\sigma}}(\mathbf{F}^{e}, T, \alpha, ^{\chi}\phi, \nabla^{\chi}\phi)$$

$$\boldsymbol{a} = \hat{\boldsymbol{a}}(\mathbf{F}^{e}, T, \alpha, ^{\chi}\phi, \nabla^{\chi}\phi), \ \mathbf{\underline{b}} = \mathbf{\underline{b}}(\mathbf{F}^{e}, T, \alpha, ^{\chi}\phi, \nabla^{\chi}\phi)$$

erive the state laws (Coleman and Noll, 1963)

$$\underline{\sigma} = \rho \frac{\partial \hat{\psi}}{\partial \underline{\mathsf{F}}^{e}} \cdot \underline{\mathsf{F}}^{eT}, \ \eta = -\frac{\partial \hat{\psi}}{\partial T}, \ X = \rho \frac{\partial \hat{\psi}}{\partial \alpha}, \quad \mathbf{a} = \frac{\partial \hat{\psi}}{\partial x \phi}, \quad \underline{\mathbf{b}} = \frac{\partial \hat{\psi}}{\partial \nabla x \phi}$$

• Residual dissipation $D^{res} = W^p - X\dot{\alpha} - \frac{\mathbf{q}}{\tau} \cdot \nabla T \ge 0$

The micromorphic approach to plasticity

D

The micromorphic approach (3)

• Take a simple quadratic potential

$$\psi(\mathbf{E}, T, \alpha, {}^{\chi}\phi, \nabla^{\chi}\phi) = \psi_1(\mathbf{E}, \alpha, T) + \psi_2(e = \phi - {}^{\chi}\phi, \nabla^{\chi}\phi, T)$$
$$\rho\psi_2 = \frac{1}{2}H_{\chi}(\phi - {}^{\chi}\phi)^2 + \frac{1}{2}A\nabla^{\chi}\phi.\nabla^{\chi}\phi$$
$$\mathbf{a} = \rho \frac{\partial\psi}{\partial\chi\phi} = -H_{\chi}(\phi - {}^{\chi}\phi), \quad \mathbf{\underline{b}} = \rho \frac{\partial\psi}{\partial\nabla\chi\phi} = A\nabla^{\chi}\phi$$

• Simple form of the partial differential equation (homogeneous, isothermal...)

$$a = \operatorname{div} \mathbf{\underline{b}} \implies {}^{\chi}\phi - \frac{A}{H_{\chi}}\Delta^{\chi}\phi = \phi$$

Helmholtz equation with a minus sign and a source term

• Coupling modulus H_{χ} and characteristic length of the medium

$$I_c^2 = \frac{A}{H_{\lambda}}$$

Stability

 $H_{\chi} > 0, \quad A > 0$

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Micromorphic continuum

Micromorphic continuum according to (Eringen and Suhubi, 1964; Mindlin, 1964)

• Select variable:

$$\phi \equiv \mathbf{F}, \quad {}^{\chi}\phi \equiv \mathbf{\chi}$$

$$p^{(i)} = \sigma : \nabla \underline{\dot{\mathbf{u}}} + \underline{\mathbf{a}} : \dot{\chi} + \underline{\underline{\mathbf{B}}} \cdot \nabla \dot{\chi}$$

 application of the principle of (infinitesimal) material frame indifference, (infinitesimal) change of observer of rate w:

$$\nabla \underline{\dot{\mathbf{u}}} \Longrightarrow \nabla \underline{\dot{\mathbf{u}}} + \mathbf{w}, \quad \chi \Longrightarrow \chi + \mathbf{w}$$

 $\Longrightarrow \sigma + {\tt a}$ must be symmetric. Rewrite the virtual power:

$$p^{(i)} = \sigma: \dot{\varepsilon} + \mathbf{s}: (\nabla \underline{\dot{\mathbf{u}}} - \dot{\chi}) + \mathbf{S}: \nabla \dot{\chi}$$

• two balance equations:

$$\operatorname{div}\left(\underline{\sigma}+\underline{s}\right)+\rho\underline{\mathbf{f}}=\mathbf{0},\quad\operatorname{div}\underline{\mathbf{S}}+s=\mathbf{0}$$

Microstrain continuum

Microstrain continuum after (Forest and Sievert, 2006)

Select

$$\phi \equiv \mathbf{\underline{C}} = \mathbf{\underline{E}}^T \cdot \mathbf{\underline{E}}, \quad {}^{\chi} \phi \equiv {}^{\chi} \mathbf{\underline{C}}, \quad \text{or} \quad \phi \equiv \boldsymbol{\underline{\varepsilon}}, \quad {}^{\chi} \phi \equiv {}^{\chi} \boldsymbol{\underline{\varepsilon}}$$
$$p^{(i)} = \boldsymbol{\underline{\sigma}} : \boldsymbol{\underline{\dot{\varepsilon}}} + \mathbf{\underline{a}} : {}^{\chi} \boldsymbol{\underline{\dot{\varepsilon}}} + \mathbf{\underline{b}} : \mathbf{\nabla}^{\chi} \boldsymbol{\underline{\varepsilon}}$$

• Constitutive coupling between macro and microstrain via the relative strain

$$\begin{split} & \underbrace{\mathbf{e}} := \boldsymbol{\varepsilon} - {}^{\boldsymbol{\chi}} \boldsymbol{\varepsilon} \\ \psi(\boldsymbol{\varepsilon}^{\boldsymbol{e}}, \quad \boldsymbol{T}, \quad \boldsymbol{\alpha}, \quad \underbrace{\mathbf{e}} := \boldsymbol{\varepsilon} - {}^{\boldsymbol{\chi}} \boldsymbol{\varepsilon}, \quad \underbrace{\mathbf{K}}_{\boldsymbol{\varkappa}} := \boldsymbol{\nabla}^{\boldsymbol{\chi}} \boldsymbol{\varepsilon}) \end{split}$$

• Take a quadratic potential

$$\mathbf{a} = H_{\chi} \mathbf{e}, \quad \mathbf{b} = A \nabla^{\chi} \mathbf{e}$$

• Extra-balance equation

$$^{\chi}\varepsilon - l_c^2 \Delta^{\chi}\varepsilon = \varepsilon, \quad ext{with} \quad l_c^2 = rac{A}{H_{\chi}}$$

example: microfoams

(Dillard et al., 2006)

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The micromorphic approach to plasticity

Cosserat continuum

Cosserat continuum

$$\phi = \mathbf{\hat{R}}, \quad {}^{\chi}\phi \equiv {}^{\chi}\mathbf{\hat{R}}$$
$$p^{(i)} = \sigma^{s}: \dot{\varepsilon} - \sigma^{a}: ((\nabla \underline{\dot{u}})^{a} - \mathbf{\dot{R}}.\mathbf{\hat{R}}^{T}) + \mathbf{\hat{M}}: \dot{\kappa}$$

name	number	DOF	DOF DOF	
	of DOF	(finite case)	(infinitesimal case)	
Cauchy	3	<u>u</u>	Ц	[Cauchy, 1823]
micromorphic	12	$\underline{u}, \chi_{\sim}$	$\underline{\mathbf{u}}, \underline{\chi}^s + \underline{\chi}^s$	[Eringen, 1964] [Mindlin, 1964]

name	number	DOF	DOF	references
	of DOF	(finite case)	(infinitesimal case)	
Cauchy	3	<u>u</u>	<u>u</u>	[Cauchy, 1823]
microdilatation	4	<u>u</u> , χ	<u>υ</u> , χ	[Goodman, Cowin, 1972] [Steeb, Diebels, 2003]
micromorphic	12	$\underline{u}, \chi_{\widetilde{\sim}}$	$\underline{\mathbf{u}}, \underline{\chi}^s + \underline{\chi}^s$	[Eringen, 1964] [Mindlin, 1964]

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Cauchy	3	<u>u</u>	Ц	[Cauchy, 1823]
microdilatation	4	<u>μ</u> , χ	<u>u</u> , <u>x</u>	[Goodman, Cowin, 1972] [Steeb, Diebels, 2003]
Cosserat	6	<u>u</u> , R	<u>и</u> , <u>Ф</u>	[Kafadar, Eringen, 1976]
micromorphic	12	$\underline{u}, {\boldsymbol{\chi}}$	$\underline{\mathbf{u}}, \underline{\chi}^s + \underline{\chi}^s$	[Eringen, 1964] [Mindlin, 1964]

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microdilatation	4	<u>υ</u> , χ	<u>u</u> , <i>x</i>	[Goodman, Cowin, 1972] [Steeb, Diebels, 2003]	
Cosserat	6	<u>u</u> , ℝ	<u>ш</u> , <u>Ф</u>	[Kafadar, Eringen, 1976]	
microstretch	7	$\underline{\mathbf{u}}, \chi, \mathbf{R}$	<u>u</u> , χ, <u>Φ</u>	[Eringen, 1990]	
micromorphic	12	$\underline{u}, \chi_{\sim}$	$\underline{u}, \chi^{s} + \chi^{a}$	[Eringen, 1964] [Mindlin, 1964]	

The micromorphic approach to plasticity

name	number	DOF DOF		references	
	of DOF	(finite case) (infinitesimal case)			
Cauchy	3	<u>u</u>	Ц	[Cauchy, 1823]	
microdilatation	4	<u>υ</u> , χ	<u>u</u> , <u>x</u>	[Goodman, Cowin, 1972] [Steeb, Diebels, 2003]	
Cosserat	6	<u>u</u> , R	<u>и</u> , <u>Ф</u>	[Kafadar, Eringen, 1976]	
microstretch	7	$\underline{\mathbf{u}}, \chi, \mathbf{R}$	<u>υ</u> , χ, <u>Φ</u>	[Eringen, 1990]	
microstrain	9	<u>u</u> , [∞] [∞] [∞]	$\underline{\mathbf{u}}, \overset{\chi}{\varepsilon}$	[Forest, Sievert, 2006]	
micromorphic	12	$\underline{u}, \chi_{\sim}$	$\underline{\mathbf{u}}, \underline{\chi}^s + \underline{\chi}^s$	[Eringen, 1964] [Mindlin, 1964]	

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Internal constraint and gradient of internal variable approach

• Impose the internal constraint that

 ${}^{\chi}\phi\simeq\phi\quad\Longrightarrow\quad\mathbf{\underline{K}}\simeq\mathbf{\nabla}\phi$

Then, the generalized stress *a* becomes a Lagrange multiplier.

- Examples
 - * $\phi \equiv \mathbf{F}$ second gradient model (Mindlin, 1965) * $\phi \equiv \mathbf{p}$ Aifantis model (Aifantis, 1987; Fleck and Hutchinson, 2001) * $\phi \equiv \varepsilon^{\mathbf{p}}$ strain gradient plasticity (Forest and Sievert, 2003; Gurtin, 2003)

$$\boldsymbol{p}^{(i)} = \boldsymbol{\sigma} : \boldsymbol{\dot{\varepsilon}}^{e} + \boldsymbol{\underline{s}} : \boldsymbol{\dot{\varepsilon}}^{p} + \boldsymbol{\underline{S}} : \boldsymbol{\dot{\underline{\varepsilon}}}^{p}, \quad \operatorname{div} \boldsymbol{\underline{S}} = \boldsymbol{\underline{s}} - \boldsymbol{\sigma}^{dev}$$

The yield condition becomes a PDE (Aifantis, Fleck-Hutchinson, Gurtin):

$$\sigma_{eq} = \sigma_{Y} + Hp - A\Delta p$$

- What do the boundary conditions become?
 - ★ For the second gradient theory, intricate b.c. involving surface curvature
 - * For gradient of plastic strain, $\underline{\underline{S}}.\underline{\underline{n}} = \underline{\underline{m}}$

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Processing of nickel foams

from PU template





electrolyte tank







Foam sheet axes and cell shape



Anisotropic Compressible Continuum Plasticity Model

Yield criterion

$$f(\underline{\sigma}) = \sigma_{eq} - R, \quad \sigma_{eq} = \left(\frac{3}{2}C\underline{\sigma}^{\text{dev}} : \underbrace{H}_{\infty} : \underline{\sigma}^{\text{dev}} + F(\underline{P} : \underline{\sigma})^2\right)^{\frac{1}{2}}$$
$$[H] = \begin{bmatrix} h_a & 0 & 0 & 0 & 0 & 0\\ 0 & h_b & 0 & 0 & 0 & 0\\ 0 & 0 & h_c & 0 & 0 & 0\\ 0 & 0 & 0 & h_d & 0 & 0\\ 0 & 0 & 0 & 0 & h_e & 0\\ 0 & 0 & 0 & 0 & 0 & h_f \end{bmatrix}, \quad [P] = \begin{bmatrix} p & 0 & 0\\ 0 & q & 0\\ 0 & 0 & r \end{bmatrix}$$

Normality rule

$$\dot{\varepsilon}^{p} = \dot{p} \frac{\partial f}{\partial \sigma} = \dot{p} \left(\frac{3}{2} C \underbrace{\mathsf{H}}_{\simeq} : \overset{\circ}{\sigma}^{\mathrm{dev}} + F(\underbrace{\mathsf{P}}_{\sim} : \overset{\circ}{\sigma}) \underbrace{\mathsf{P}}_{\sim} \right)$$

Consistency condition (cumulative plastic strain p)

$$\dot{p} = \frac{\frac{\partial f}{\partial \sigma} : \mathbf{E} : \dot{\mathbf{E}}}{\frac{\partial f}{\partial \sigma} : \mathbf{E} : \dot{\mathbf{E}} : \frac{\partial f}{\partial \sigma} + \frac{\partial R}{\partial p}}$$

Nonlinear isotropic hardening

$$R = R_0 + Hp + Q(1 - \exp(-bp))$$

Application to nickel foams

Parameter identification

C	F	R ₀ (MPa)	Q (MPa)	b	H (MPa)	h _a	h _b	h _c	h _d	h _e	h _f
1	6,8.10-4	0,30	0,37	3,3	8,17	0,41	1,8	1,75	1,45	1	1



р	q	r
1	37	23

Metal foam as a micromorphic continuum : The microfoam model

Degrees of freedom at material point

displacement vector $\underline{\mathbf{u}}$, microdeformation tensor $\boldsymbol{\chi}$

after [Mindlin, 1964] [Eringen, 1964] Deformation measures

strain tensor $\boldsymbol{\varepsilon} = (\underline{\mathbf{u}} \otimes \boldsymbol{\nabla} + \boldsymbol{\nabla} \otimes \underline{\mathbf{u}})/2$

relative deformation tensor $\mathbf{e} = \mathbf{\underline{u}} \otimes \boldsymbol{\nabla} - \boldsymbol{\chi}$

microdeformation gradient $\mathbf{K} = \chi \otimes \nabla$

Power density of internal forces / generalized stress tensors

$$p^{(i)} = \sigma : \dot{\varepsilon} + \underline{s} : \dot{\underline{e}} + \underline{\underline{M}} : \dot{\underline{K}}$$

Balance of momentum

Balance of micromorphic momentum

Boundary conditions (traction and double traction vectors)

$$\underline{\mathbf{t}} = (\underline{\sigma} + \underline{\mathbf{s}}) \cdot \underline{\mathbf{n}}, \quad \underline{\mathsf{T}} = \underline{\underline{\mathsf{M}}} \cdot \underline{\mathbf{n}}$$

Variational formulation for finite element implementation

$$\int_{V} p^{(i)} dV = \int_{\partial V} \left(\underline{\mathbf{t}} \cdot \underline{\dot{\mathbf{n}}} + \underline{\mathbf{T}} : \dot{\mathbf{\chi}} \right) dS$$

Application to nickel foams

 $(\underbrace{\boldsymbol{\sigma}}_{\widetilde{\boldsymbol{\omega}}} + \underbrace{\mathbf{s}}_{\widetilde{\boldsymbol{\omega}}}) \cdot \boldsymbol{\nabla} = \mathbf{0}$ $\underbrace{\mathbf{M}}_{\widetilde{\boldsymbol{\omega}}} \cdot \boldsymbol{\nabla} + \underbrace{\mathbf{s}}_{\widetilde{\boldsymbol{\omega}}} = \mathbf{0}$

The microfoam model : Constitutive equations

Strain partition

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}^{e} + \boldsymbol{\varepsilon}^{p}, \quad \boldsymbol{e} = \boldsymbol{e}^{e} + \boldsymbol{e}^{p}, \quad \boldsymbol{K} = \boldsymbol{K}^{e} + \boldsymbol{K}^{p}$$

Linear elasticity

simplified moduli as proposed by [Shu et al., 1999] $\underline{\mathbf{M}} = l_c^2 \underbrace{\mathbf{c}}_{\simeq} : \underbrace{\mathbf{K}}^e$ characteristic length l_c when $||\underline{\mathbf{a}}||$ is very large, and if $\underline{\mathbf{e}}^p = 0$, then $\underline{\mathbf{u}} \otimes \boldsymbol{\nabla} \simeq \boldsymbol{\chi}$ (second gradient theory)

The microfoam model : Constitutive equations

Yield criterion

$$f(\underline{\sigma}) = \sigma_{eq} - R$$

$$\sigma_{eq} = \left(\frac{3}{2}C\underline{\sigma}^{\text{dev}} : \underbrace{\mathbf{H}}_{\approx} : \underline{\sigma}^{\text{dev}} + F(\underbrace{\mathbf{P}}_{\approx} : \underline{\sigma})^2 + a_1\underline{\mathbf{s}}^{\text{dev}} : \underline{\mathbf{s}}^{\text{dev}} + a_2\underline{\mathbf{s}}^{\text{dev}} : \underline{\mathbf{s}}^{\text{devT}} - b_1\underline{\mathbf{M}} : \underline{\mathbf{M}}\right)^{\frac{1}{2}}$$

Normality rule

$$\dot{\varepsilon}^{p} = \dot{p} \frac{\partial f}{\partial \sigma}, \quad \dot{\mathbf{e}}^{p} = \dot{p} \frac{\partial f}{\partial \mathbf{s}}, \quad \dot{\mathbf{K}}^{p} = \dot{p} \frac{\partial f}{\partial \mathbf{M}}$$

Microfoam: $a_1 = a_2 = a_3 = b_1 = 0$, $||\mathbf{a}||$ very large constrained microdeformation (second gradient), linear relationship between gradient of microdeformation and hyperstress tensor

Microstrain continuum

degrees of freedom $(\underline{\mathbf{u}}, \underline{\chi}^s)$ strain measures : $(\underline{\varepsilon}, \underline{\varepsilon} - \underline{\chi}^s, \underline{\chi}^s \otimes \nabla)$

$$egin{aligned} & \mathbf{g} = \mathbf{C} : (\mathbf{g} - \mathbf{g}^{p}) \ & \mathbf{s} = b(\mathbf{g} - \mathbf{\chi}^{s}) \ & \mathbf{S} = A \mathbf{\chi}^{s} \otimes \mathbf{\nabla} \end{aligned}$$

$$\operatorname{div}\left(\underline{\sigma}+\underline{\mathbf{s}}\right)=\mathbf{0}$$

$$\operatorname{div} \underline{\underline{S}} + \underline{\underline{s}} = 0$$

$$egin{aligned} S_{ijk,k}+s_{ij}&=0 \ A\chi^s_{ij,kk}+b(arepsilon_{ij}-\chi^s_{ij})&=0 \end{aligned}$$

$$\varepsilon_{ij} = \chi^s_{ij} - l^2 \Delta \chi^s_{ij}$$

Link with "implicit gradient–enhanced elastoplasticity models" [Engelen *et al.*, 2003]

Boundary conditions : (u_i, χ_{ij}^s) or $(\sigma_{ij} + s_{ij})n_j, S_{ijk}n_k$

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Fracture of nickel foams



simple fracture criterion $p = p_{crit} = 0.08$ limited scatter in p_{crit}

softening law for $p > p_{crit}$

$$R = R(p > p_{crit})$$

tensile test in direction RD

Fracture anisotropy of nickel foams

prediction of fracture stress for tension along TD



Tension of central crack nickel foam plate



crack length : 10 mm

Application to nickel foams

Simulation with a classical compressible plasticity model





Simulation with a classical compressible plasticity model





Simulation with a classical compressible plasticity model





Simulation with the micromorphic foam model



Damage zone size



Convergence of the load-displacement curve



Application to nickel foams

Fracture anisotropy of nickel foams

prediction of fracture stress for tension along TD



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Behaviour of Nickel foils



elastoplastic model with linear isotropic hardening



Micromechanical approach





















Mesh sensitivity



Prediction of elastoplastic properties



Eight cell volume element



3 600 000 DOF - 36 GB - 300 hours

Micromechanical approach

Eight cell volume element



3 600 000 DOF - 36 GB - 300 hours - red: p > 0.03

Eight cell volume element



3 600 000 DOF - 36 GB - 300 hours - red: p > 0.03

Eight cell vs. one cell volume element



Micromechanical approach

Perspectives

- Micromorphic approach in the mechanics of generalized continua
- Plasticity and fracture with the *microfoam* model
- Micromechanical definition of the additional degrees of freedom :

link between the microdeformation tensor and the deformation of individual cells

 \implies see the presentation of Anthony Burtheau



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