

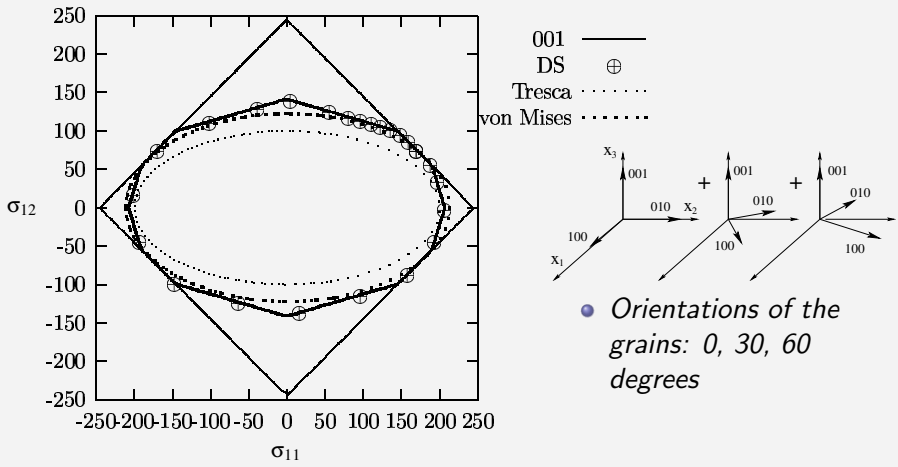
Contents

From single crystal to polycrystal
 Initial yield surfaces for DS material
 Initial yield surfaces for a polycrystalline material

Scale transition rules
 Various rules and their physical meaning

Main features of the polycrystal models
 Subsequent yield surfaces
 Complex paths
 Examples of FE computations using crystal plasticity
 Twinning and phase transformation

Yield surface for DS material with 3 orientations



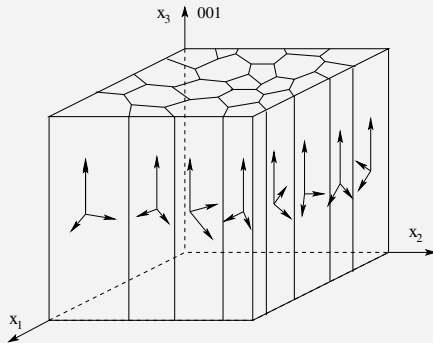
More slip directions in the plane $\sigma_{11}-\sigma_{12}$: the model turns out to become similar to von Mises/Tresca

Use of uniform field models in Crystal plasticity

Georges Cailletaud

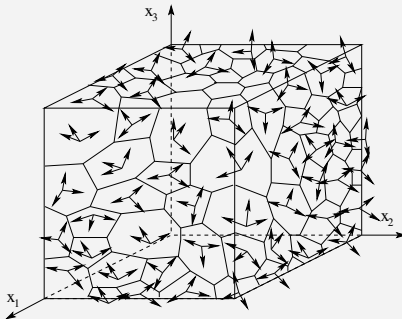
*Centre des Matériaux
 Ecole des Mines de Paris/CNRS*

Loading surface of a directionally solidified material (DS)



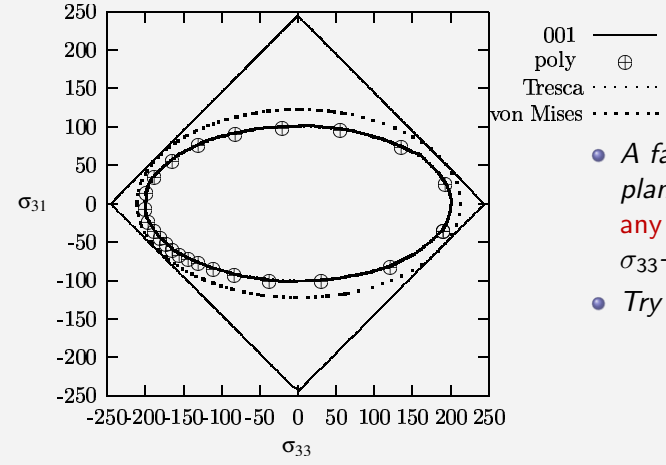
- For a DS material, all the grains have a common crystal axis, say (001). The orientation of each grain is then defined by one angle around this axis.
- Assuming that elasticity is uniform, the stress is also uniform during the elastic phase.

Loading surface of a polycrystalline material



- The orientation of a grain is defined by a set of three Euler angles, (ϕ_1, Φ, ϕ_2)
- Assume a uniform elasticity

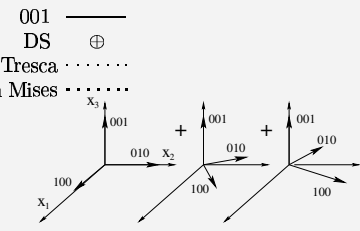
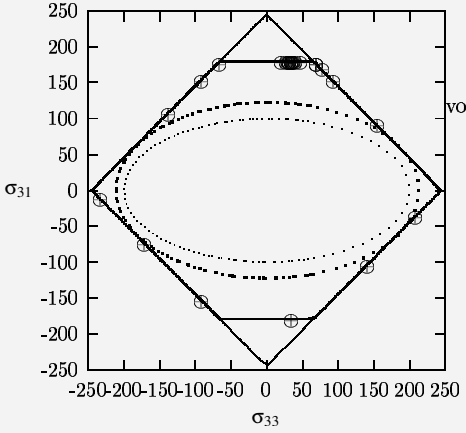
Yield surface for a polycrystalline aggregate with 100 orientations



- A favourably oriented plane can be found for any direction in the $\sigma_{33}-\sigma_{31}$ plane
- Try an other set [HERE](#)

Tresca like criterion is obtained for plane $\sigma_{33}-\sigma_{31}$

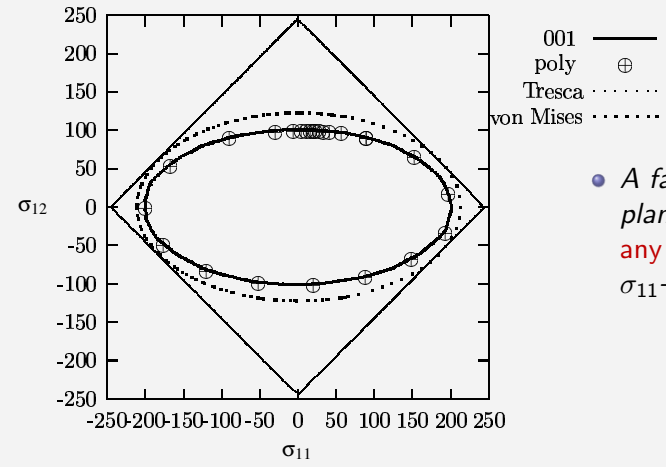
Yield surface for DS material with 3 orientations



- Orientations of the grains: 0, 30, 60 degrees
- Try an other set [HERE](#)

No additional slip activated by shear in the plane $\sigma_{33}-\sigma_{13}$: the yield stress in shear remains high, and sharp corners are still present

Yield surface for a polycrystalline aggregate with 100 orientations



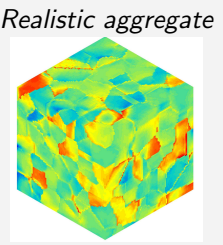
- A favourably oriented plane can be found for any direction in the $\sigma_{11}-\sigma_{12}$ plane

Tresca like criterion is obtained for plane $\sigma_{11}-\sigma_{12}$

Computations at various scales



Level (1)
Macroscopic models
Average stress and strain tensor



Level (3)
Local information
Respect local Constitutive Equations
Equilibrium
Lectures 4 & 5

Intermediate level
Averaging process
(2) phase by phase
(2') grain by grain

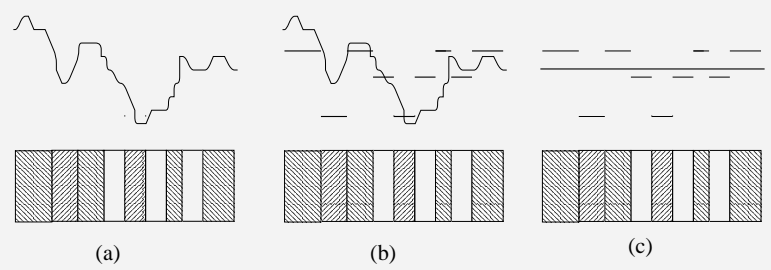
Level (2) or (2')
Local average
Respect local Constitutive Equations
no neighboring effect
This lecture

Rules used for the scale transition in uniform field models

- Static, uniform stress, $\bar{\sigma}^g = \sigma$
- From Taylor to Kröner
 - Taylor [Taylor, 1938], uniform plastic strain, $\bar{\epsilon}^{Pg} = \epsilon^P$
 - Lin-Taylor [Lin, 1957], uniform total strain, $\bar{\epsilon}^g = \epsilon$
 - Kröner [Kröner, 1971], elastic accommodation $\bar{\sigma}^g = \sigma + \mu(\epsilon^P - \epsilon^{Pg})$
- Tangent and secant approximations
 - Hill [Hill, 1965], elastoplastic accommodation $\dot{\bar{\sigma}}^g = \dot{\sigma} + \underline{\underline{L}}^* : (\dot{\epsilon} - \dot{\epsilon}^g)$
 - Berveiller-Zaoui [Berveiller and Zaoui, 1979] estimation, with α varying typically between 1 and 0.001 (isotropic elasticity, spherical inclusions) $\bar{\sigma}^g = \sigma + \mu\alpha(\epsilon^P - \epsilon^{Pg})$
- Viscous and viscoplastic scheme
 - Budianski, Hutchinson, Molinari, ... [Hutchinson, 1966a, Molinari et al., 1987a]
 - Translated fields [Sabar et al., 2002], $\dot{\bar{\sigma}}^g = \dot{\sigma} + 2\mu(1 - \beta)(\frac{5\eta}{3\eta + 2\eta^g}\dot{\epsilon}^v - \dot{\epsilon}^{vg})$
- Parametric scale transition rule
 - Cailletaud, Pilvin, β -model

Various scales in heterogeneous material modeling

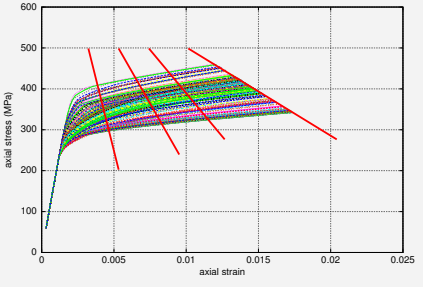
- (a) Microstructure calculations: take into account local phases and local equilibrium (Next two lectures)
- (b) Uniform field models: take into account local phases only (This lecture)
- (c) Macroscopic models: do not account for local phases nor local equilibrium



Definitions used in the transition rules

- Phase characteristics
 - Chemical composition, shape, crystallographic orientation
 - f^g is the volume fraction of the phase g
 - The choice of the representation of a phase is driven by the contrast between properties. Two phase polycrystal \rightarrow two phases, discriminating by the chemical nature, or the crystal network; one phase polycrystal \rightarrow N crystallographic phases
- Notations:
 - Stress in the phase g , macroscopic stress: $\bar{\sigma}^g, \bar{\sigma} = \sum_g f^g \bar{\sigma}^g$
 - Strain in the phase g , macroscopic strain: $\bar{\epsilon}^g, \bar{\epsilon} = \sum_g f^g \bar{\epsilon}^g$
 - For an uniform local elasticity, the macroscopic plastic strain is also the average of the local plastic strains: $\bar{\epsilon}^P = \sum_g f^g \bar{\epsilon}^{Pg}$

Physical meaning of Berveiller–Zaoui rule



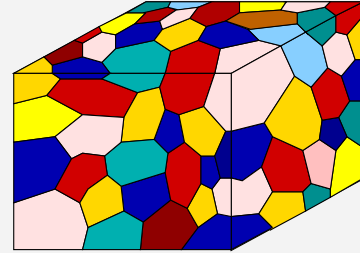
- Uniform elasticity
- Incompressible plastic flow
- Radial loading path
- General expression, using μ' , actual shear modulus, β' = $2(4 - 5\nu')/15(1 - \nu')$

$$\dot{\underline{\sigma}}^g = \dot{\underline{\sigma}} + 2\mu \frac{\mu'(1 - \beta')}{\beta'\mu + (1 - \beta')\mu'} (\dot{\underline{\varepsilon}}^P - \dot{\underline{\varepsilon}}^{Pg})$$

- Pure tension, assuming $\nu = 1/2$ and introducing $H = \sigma/\varepsilon^P$:
$$\dot{\sigma}_g = \dot{\sigma} + \frac{\mu H}{H + 2\mu} (\dot{\varepsilon}^P - \dot{\varepsilon}^{Pg})$$
- Onset of plastic flow: Kröner's rule
- The accommodation factor $C = (\sigma^g - \sigma)/(\varepsilon^P - \varepsilon^{Pg})$ decreases when plastic strain increases

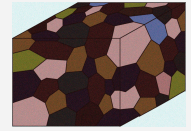
Calibration of the scale transition rule

- FE computation



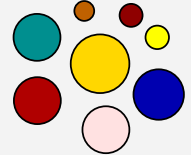
- 1 Perform a FE computation of a realistic aggregate
- 2 Post-treatment to get the macroscopic stress and strain
- 3 Post-treatment to get the average values in each phase
- 4 Identification of the parameters of the transition rule to have a good fit for both global and local levels

- Global response



$\underline{\sigma}$
 $\underline{\varepsilon}$

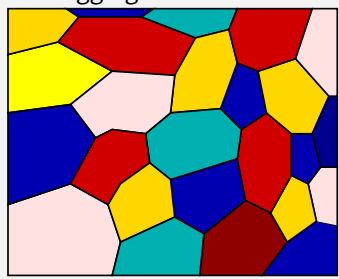
- Local response



$\underline{\sigma}^g$
 $\underline{\varepsilon}^g$

Hill's self-consistent model

- Real aggregate



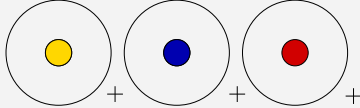
$$\dot{\underline{\sigma}} = \underline{\underline{\mathbb{L}}}^* : (\dot{\underline{\varepsilon}} - \dot{\underline{\varepsilon}})$$

with $\underline{\underline{\mathbb{L}}}^*$, accommodation tensor:

$$\underline{\underline{\mathbb{L}}}^* = \underline{\underline{\mathbb{L}}}^{\text{eff}} : (\underline{\underline{\mathbb{S}}}^{-1} - \underline{\underline{\mathbb{I}}})$$

$$(\underline{\underline{\mathbb{L}}} + \underline{\underline{\mathbb{L}}}^*) : \dot{\underline{\varepsilon}} = (\underline{\underline{\mathbb{L}}}^{\text{eff}} + \underline{\underline{\mathbb{L}}}^*) : \dot{\underline{\varepsilon}}$$

- Modelled as a collection of auxiliary problems



$$\dot{\underline{\sigma}} = \underline{\underline{\mathbb{L}}} : (\underline{\underline{\mathbb{L}}} + \underline{\underline{\mathbb{L}}}^*)^{-1} : (\underline{\underline{\mathbb{L}}}^{\text{eff}} + \underline{\underline{\mathbb{L}}}^*) : \dot{\underline{\varepsilon}}$$

- implicit problem, since the HEM is not known

So that :

$$\underline{\underline{\mathbb{L}}}^{\text{eff}} = \langle \underline{\underline{\mathbb{L}}} : (\underline{\underline{\mathbb{L}}} + \underline{\underline{\mathbb{L}}}^*)^{-1} : (\underline{\underline{\mathbb{L}}}^{\text{eff}} + \underline{\underline{\mathbb{L}}}^*) \rangle$$

Introduction of a parametric scale transition rule

- The local stress decrease when the grain becomes more plastic than the matrix

$$\underline{\underline{\sigma}}^g = \underline{\underline{\sigma}} + C (\underline{\underline{\mathbb{B}}} - \underline{\underline{\beta}}^g)$$

$$\underline{\underline{\mathbb{B}}} = \sum_g f_g \underline{\underline{\beta}}^g$$

f_g is the volume fraction of phase g , $\underline{\underline{\beta}}^g$ characterizes the state of redistribution

- New interphase accommodation variables $\underline{\underline{\beta}}^g$, with a kinematic evolution rule [Cailletaud 87]

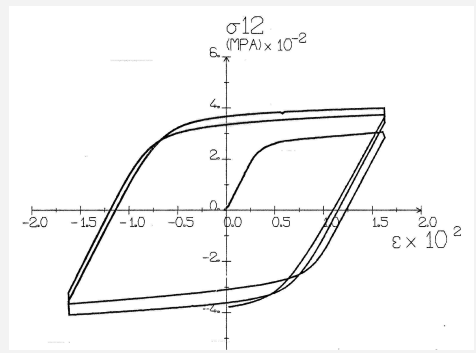
- rule 1: $\dot{\underline{\underline{\beta}}}^g = \dot{\underline{\underline{\varepsilon}}}^g - D \dot{\underline{\underline{\varepsilon}}}_{\text{eq}}^g \underline{\underline{\beta}}^g$

- rule 2: $\dot{\underline{\underline{\beta}}}^g = \dot{\underline{\underline{\varepsilon}}}^g - D \dot{\underline{\underline{\varepsilon}}}_{\text{eq}}^g (\underline{\underline{\beta}}^g - \delta \underline{\underline{\varepsilon}}^g)$

- Identification procedure, inverse approach from finite element [Pilvin 97]: the coefficient C , D , δ , are not exactly material coefficients, but scale transition parameters, which should be fitted from Finite Element computation on realistic polycrystalline aggregates.

Distorsion of a yield surface

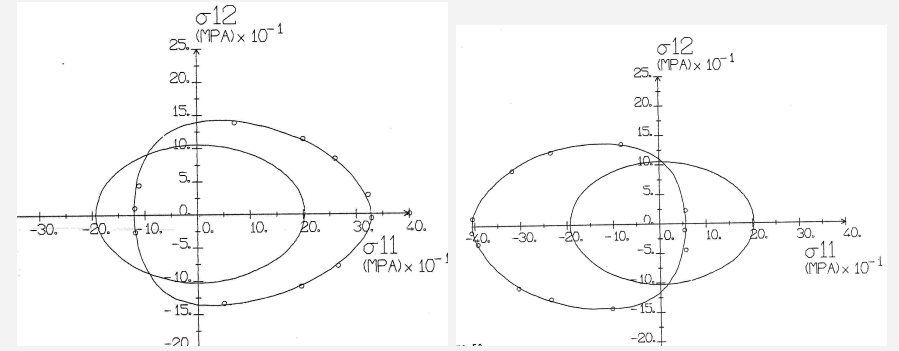
Fatigue tests on a 2024 aluminium alloy



PhD Marc Rousset, 1985

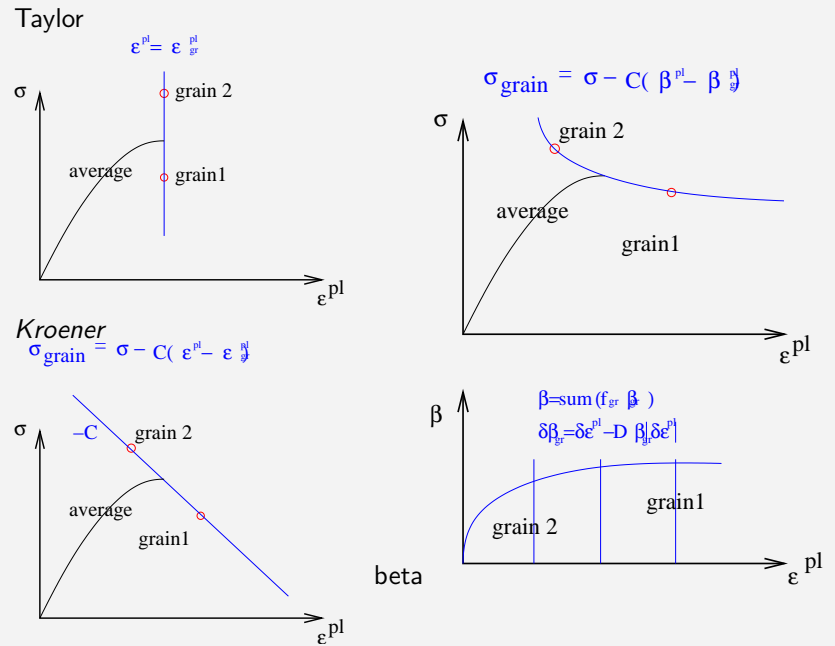
Experimental distortion of a yield surface (1bis)

Second cycle



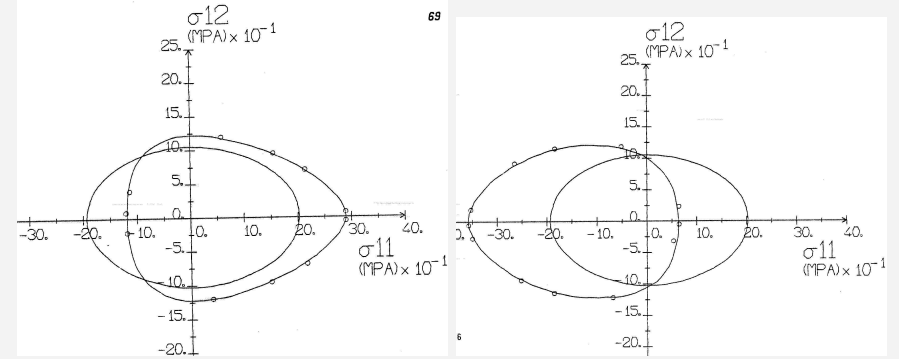
PhD Marc Rousset, 1985

Summary of the scale transition rules



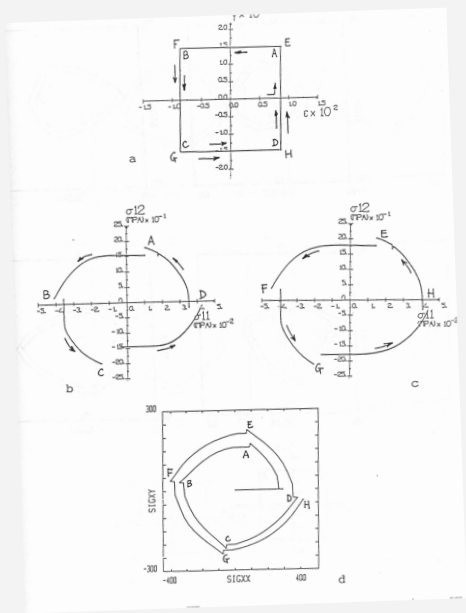
Experimental distortion of a yield surface (1)

First cycle



PhD Marc Rousset, 1985

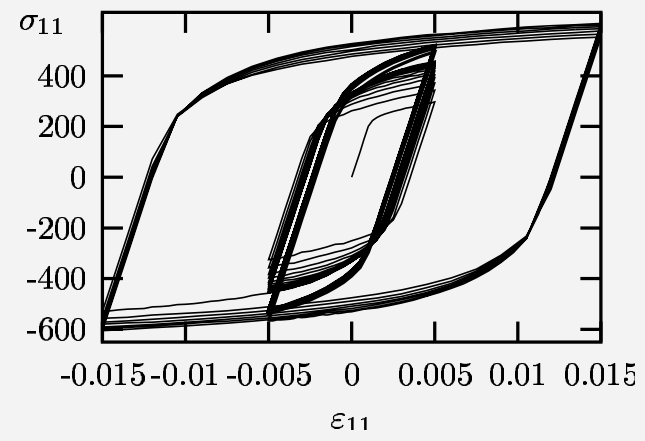
Experimental distortion of a yield surface (2)



- Square path in the $\epsilon_{11}-\epsilon_{12}$ plane
- Yield surface investigation for each corner, at 1st and 2nd cycle
- Low viscous effect
- Good prediction of the stress response in the $\sigma_{11}-\sigma_{12}$ plane

PhD Marc Rousset, 1985

Simulation of the memory effect

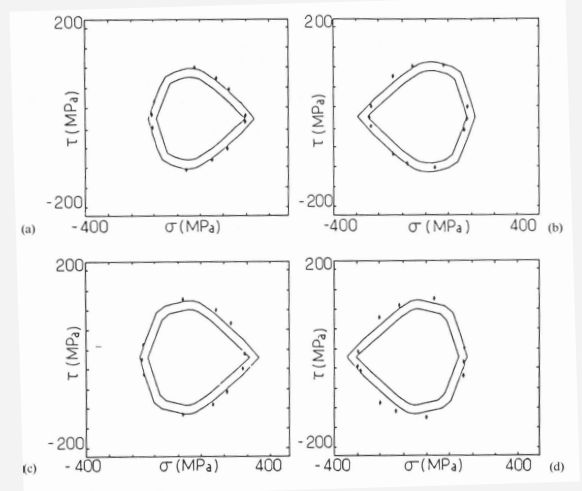


R_0	Q	b	c	d	h_i
100	50	10	5000	100	1, i

[Besson et al., 2001]

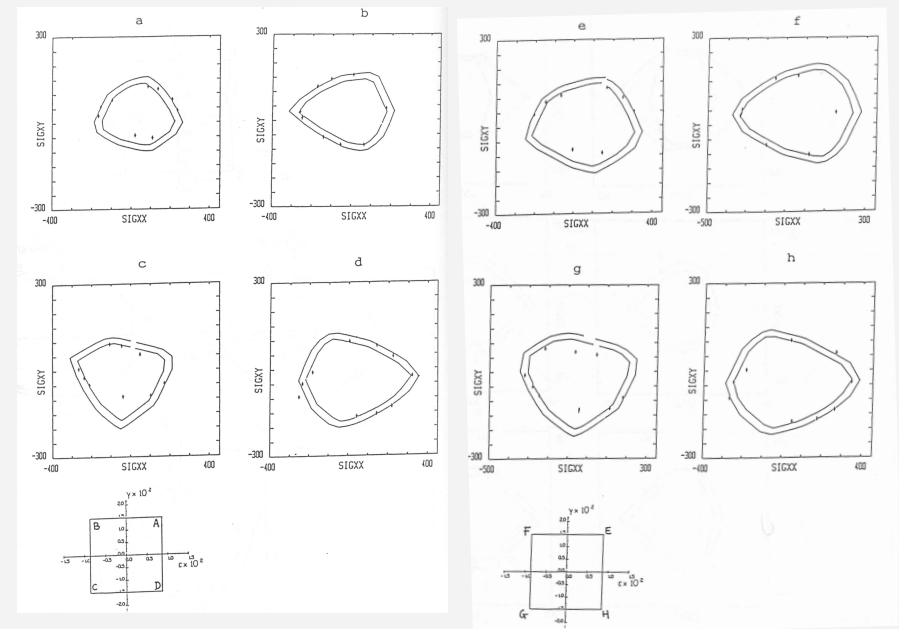
Simulation of the distortion of a yield surface

First cycle and second cycles, 2024 aluminium alloy



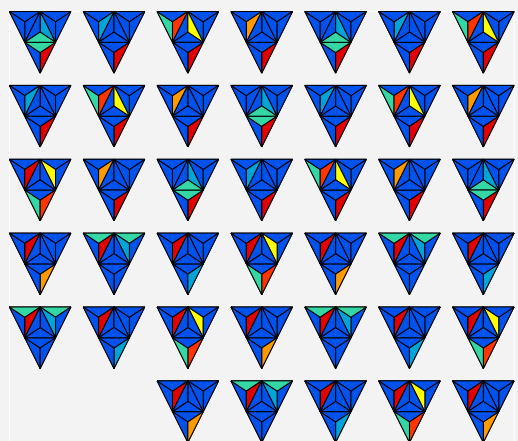
[Cailletaud, 1992]

Simulation of the distortion of a yield surface



[Cailletaud, 1992]

Contour of accumulated slip after a tension test

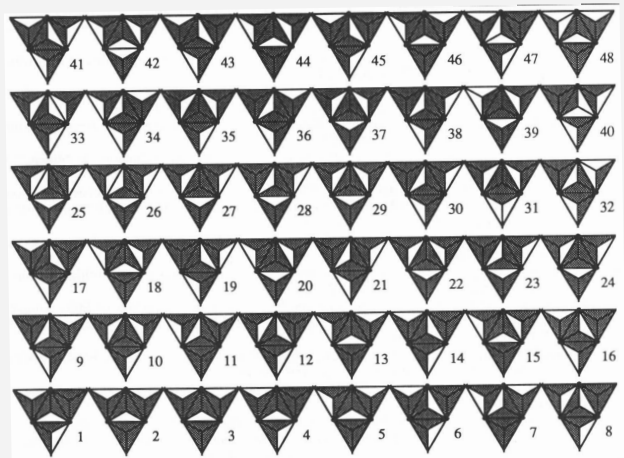


FCC grains, 12 slip systems are reported in each big triangles



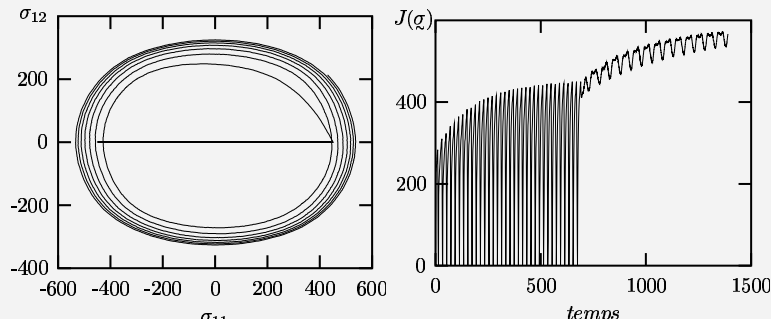
Active slip systems for an out-of-phase test

Average value per grain: 3.25



Simulation of the additional hardening obtained in out-of-phase tension-shear loading

The loading is a circle in $\epsilon_{11}-\epsilon_{12}$ plane



Response in the stress plane

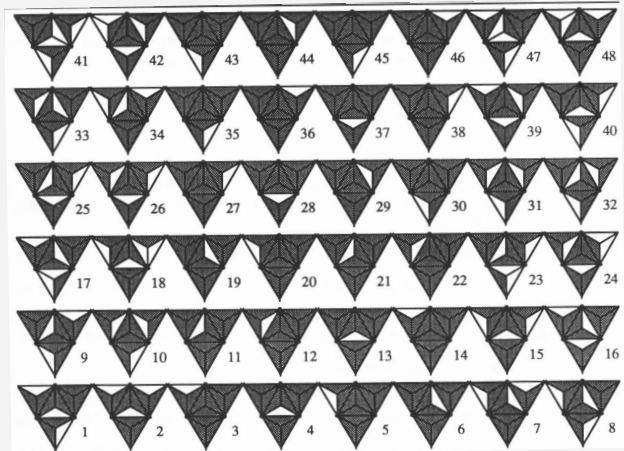
stress

Evolution of the equivalent

[Besson et al., 2001]

Active slip systems in tension

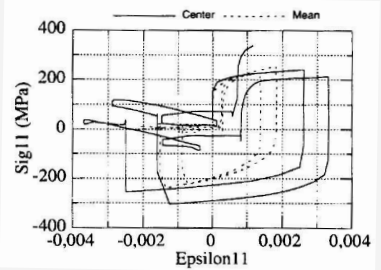
Average value per grain: 2.17



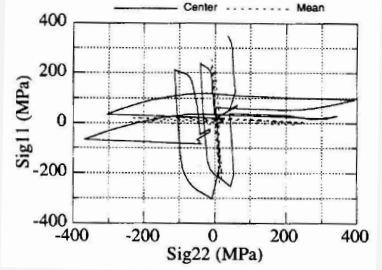
Active systems in white

A 3D specimen: computation results

A polycrystal models in each Gauss point

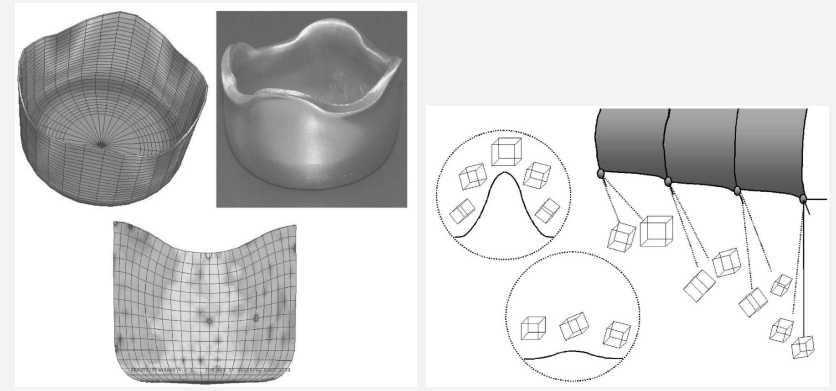


Hysteresis loop for 11 component



Loading path in the $\sigma_{11} - \sigma_{22}$ plane

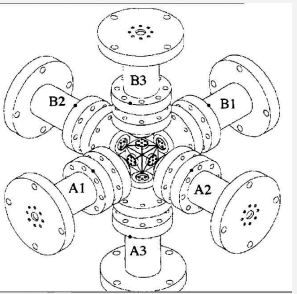
Deep drawing, one polycrystal for each Gauss point



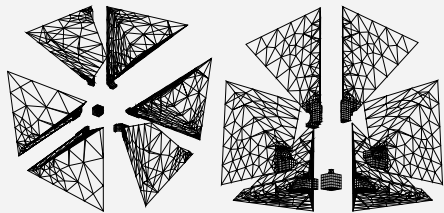
(more in [Raabe and Roters, 2004])

A 3D specimen (exp. made at LMT Cachan)

A polycrystal models in each Gauss point



Schematic view of the specimen



The mesh with subdomains for parallel computation

(more in [Feyel et al., 1997])

Indentation, one polycrystal for each Gauss point

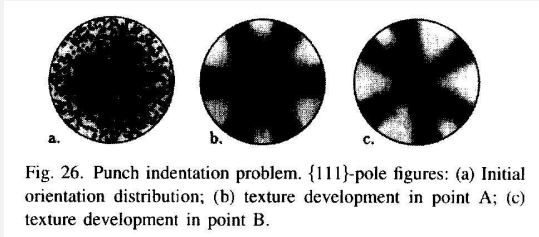
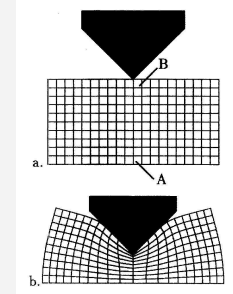


Fig. 26. Punch indentation problem. {111}-pole figures: (a) Initial orientation distribution; (b) texture development in point A; (c) texture development in point B.




(more in [Miehe et al., 1999])

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- *Polycrystal effect* [Cailletaud, 1992, Sarma and Dawson, 1996, Dawson and Marin, 1998, Nemat-Nasser and Li, 1998]
- *Scale transition rules* [Hutchinson, 1966b, Berveiller and Zaoui, 1979, Molinari et al., 1987b, Pilvin and Cailletaud, 1990, Cailletaud and Pilvin, 1994, Forest and Pilvin, 1995, Sabar et al., 2002]
- *Twining and phase transformation*
[Cherkaoui et al., 2000, Staroselski and Anand, 1998]
- *Finite Element computations with polycrystals,*
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